UNIVERSITY OF CALIFORNIA, DAVIS

Department of Electrical and Computer Engineering

EEC250

Linear Systems and Signals

Fall 2009

Lecture 7

Topics:

- a) Properties of the matrix exponential
- b) Eigenvalue/eigenvector expressions for e^{At}

Matrix exponential: Let A be an $n \times n$ real matrix. Then the matrix exponential $\Phi(t) = \exp(At)$ is defined as the solution of the differential equation

$$\Phi(t) = A\Phi(t) \tag{1}$$

for $t \geq 0$, with initial condition $\Phi(0) = I_n$. If F(s) denotes the Laplace transform of $\Phi(t)$, by Laplace transforming equation (1), we obtain

$$sF(s) - \Phi(0) = AF(s),$$

so that

$$(sI - A)F(s) = \Phi(0) = I_n.$$

This implies $F(s) = (sI - A)^{-1}$, so that $\exp(At)$ can also be obtained through the Laplace transform relation

$$\exp(At) \stackrel{LT}{\longleftrightarrow} (sI - A)^{-1} \,. \tag{2}$$

Property 1: $\exp(At)$ admits the power series expansion

$$\exp(At) = I_n + At + \frac{(At)^2}{2!} + \dots + \frac{(At)^k}{k!} + \dots,$$
(3)

for $t \geq 0$, which can be viewed as a matrix version of the power series expansion

$$\exp(at) = \sum_{k=0}^{\infty} \frac{(at)^k}{k!}$$

of the scalar exponential function. To verify (3), denote the power series on the right hand side of (3) by $\Phi(t)$ and assume that it converges and can be differentiated term by term. Differentiating each term gives

$$\frac{d}{dt}\Phi(t) = A\left(I_n + At + \dots + \frac{(At)^{k-1}}{(k-1)!} + \dots\right) = A\Phi(t)$$

for $t \geq 0$, with $\Phi(0) = I_n$, so that $\Phi(t)$ obeys the differential equation (1) defining the matrix exponential.

Another way of deriving (3) relies on the observation that for s sufficiently large, $(sI - A)^{-1}$ admits the power series expansion

$$(sI - A)^{-1} = \frac{1}{s} \left(I - \frac{A}{s} \right)^{-1} = \frac{1}{s} \left[I + \frac{A}{s} + \left(\frac{A}{s} \right)^2 + \dots + \left(\frac{A}{s} \right)^k + \dots \right]. \tag{4}$$

Then, using the fact that

$$\frac{t^k}{k!} \stackrel{\text{LT}}{\longleftrightarrow} \frac{1}{s^{k+1}}$$
,

and taking the inverse Laplace transform of (4) gives (3).

Examples: (i) Consider the $r \times r$ nilpotent matrix

$$N = \left[\begin{array}{cccc} 0 & 1 & 0 & \dots & 0 \\ & 0 & 1 & \ddots & 0 \\ & & \ddots & \ddots & \\ & & & 0 & 1 \\ 0 & & & & 0 \end{array} \right]$$

which has ones on its first superdiagonal, and zeros everywhere else. It has property that for $\ell < r$ the matrix

has ones on its ℓ -th superdiagonal, and zeros everywhere else, and $N^{\ell} = 0$ for $\ell \geq r$. Then in the power series expansion (3) for $\exp(Nt)$, only the first r terms are nonzero, so that

(ii) Let

$$A = \left[\begin{array}{cc} 0 & -1 \\ 1 & 0 \end{array} \right] .$$

Then $A^{2p} = (-1)^p I_2$ and $A^{2p+1} = (-1)^p A$, where I_2 denotes the 2 × 2 identity matrix, so that

$$\exp(At) = \left(1 - \frac{t^2}{2!} + \frac{t^4}{4!} \dots\right) I_2 + \left(t - \frac{t^3}{3!} + \frac{t^5}{5!} + \dots\right) A$$

$$= \cos t I_2 + \sin t A = \begin{bmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{bmatrix}.$$

Property 2: The matrix exponential has the transition property

$$\exp(A(t_1 + t_2)) = \exp(At_1) \exp(At_2) \tag{5}$$

for $t_1, t_2 \geq 0$. To see how this propoerty arises, consider solving the differential equation (1) for $0 \leq t \leq t_1 + t_2$. We can solve the equation in one step over the whole interval, in which case the solution at $t = t_1 + t_2$ is $\exp(A(t_1 + t_2))$. Alternatively, we can first solve the equation over $[0, t_2]$ and then over $[t_2, t_1 + t_2]$. In this case the solution at $t = t_2$ is $\exp(At_2)$, and the solution for $t_2 \leq t \leq t_1 + t_2$ is obtained by solving

$$\frac{d}{dt}\Phi(t) = A\Phi(t)$$

over $[t_2, t_1 + t_2]$ with initial condition $\Phi(t_2) = \exp(At_2)$. But the system is LTI, so that the solution at $t = t' + t_2$ with $t' \ge 0$ is given by

$$\Phi(t) = \exp(At')\Phi(t_2) .$$

Setting $t = t_1 + t_2$ and $t' = t_1$ in this identity gives (5).

Property 3: In general if A and B are two arbitrary matrices

$$\exp((A+B)t) \neq \exp(At) \exp(Bt). \tag{6}$$

To see this, let

$$A = \begin{bmatrix} 0 & -1 \\ 0 & 0 \end{bmatrix}$$
 and $B = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$.

Both A and B are nilpotent with

$$\exp(At) = \begin{bmatrix} 1 & -t \\ 0 & 1 \end{bmatrix}$$
 and $\exp(Bt) = \begin{bmatrix} 1 & 0 \\ t & 1 \end{bmatrix}$,

so

$$\exp(At)\exp(Bt) = \begin{bmatrix} 1 & -t \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ t & 1 \end{bmatrix} = \begin{bmatrix} 1-t^2 & -t \\ t & 1 \end{bmatrix}.$$

On the other hand

$$A+B=\left[\begin{array}{cc} 0 & -1\\ 1 & 0 \end{array}\right]\;,$$

and it was shown earlier that

$$\exp((A+B)t) = \begin{bmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{bmatrix}$$

which is clearly different from $\exp(At) \exp(Bt)$.

However if A and B commute, i.e., AB = BA, we have

$$\exp((A+B)t) = \exp(At) \exp(Bt). \tag{7}$$

To prove this, note that when A and B commute

$$(A+B)^k = A^k + \binom{k}{1} A^{k-1} B + \dots + \binom{k}{1} A B^{k-1} + B^k.$$

Taking this identity into account, and multiplying power series term by term gives

$$\exp(At) \exp(Bt) = \left[I + At + \frac{(At)^2}{2!} + \dots \right] \left[I + Bt + \frac{(Bt)^2}{2!} + \dots \right]$$
$$= I + (A+B)t + \frac{((A+B)t)^2}{2!} + \dots = \exp((A+B)t).$$

Examples: (i) Consider the $r \times r$ Jordan block

$$J = \left[egin{array}{cccc} \lambda & 1 & & & 0 \\ & \lambda & \ddots & & \\ & & \ddots & \ddots & \\ & & & \ddots & 1 \\ 0 & & & \lambda \end{array}
ight] = \lambda I_r + N$$

where λI_r and N commute. We have

$$e^{\lambda I_r t} = \left(1 + \lambda t + \frac{(\lambda t)^2}{2!} + \ldots\right) I_r = e^{\lambda t} I_r$$

and

$$\exp(Jt) = \exp((\lambda I_r + N)t) = \exp(\lambda I_r t) \exp(Nt)$$

$$= \exp(\lambda t) \begin{bmatrix} 1 & t & t^{r-1}/(r-1)! \\ 1 & \ddots & \\ & \ddots & \ddots \\ & & 1 & t \\ 0 & & 1 \end{bmatrix}.$$

(ii) Let

$$A = \left[\begin{array}{cc} \sigma & -\omega \\ \omega & \sigma \end{array} \right] = \sigma_2 + \omega Q$$

with

$$Q \stackrel{\triangle}{=} \left[egin{array}{cc} 0 & -1 \ 1 & 0 \end{array}
ight] \, ,$$

where I_2 and Q commute. Then

$$\exp(At) = \exp(\sigma t) \exp(\omega Qt) = \exp(\sigma t) \begin{bmatrix} \cos \omega t & -\sin \omega t \\ \sin \omega t & \cos \omega t \end{bmatrix}.$$

Property 4: $\exp(At)$ is an *invertible matrix* for all $t \geq 0$ and

$$\left(\exp(At)\right)^{-1} = \exp((-A)t). \tag{8}$$

To see this, note that A and -A commute and A + (-A) = 0, so that

$$e^{0t} = I_n = \exp(At) \exp((-A)t) .$$

Eigenvalue/eigenvector expressions: The matrix exponential $\exp(At)$ can also be expressed in terms of the eigenvalues and eigenvectors (or generalized eigenvectors) of A.

Case 1: A is diagonalizable: In this case A admits n independent eigenvectors \mathbf{p}_i corresponding to eigenvalues λ_i , i.e.

$$A\mathbf{p}_i = \lambda_i \mathbf{p}_i \tag{9}$$

for $1 \le i \le n$. The relations (9) can be combined as a single matrix equation

$$AP = P\Lambda$$

with

$$P = [\mathbf{p}_1 \dots \mathbf{p}_i \dots \mathbf{p}_n]$$
 and $\Lambda = \operatorname{diag} \{\lambda_i, 1 \leq i \leq n\}$,

so that $A = P\Lambda P^{-1}$. The matrix

$$P^{-1} = Q = \begin{bmatrix} \mathbf{q}_1^T \\ \vdots \\ \mathbf{q}_i^T \\ \vdots \\ \mathbf{q}_n^T \end{bmatrix}$$

yields the left eigenvectors of A since $P^{-1}A = P^{-1}\Lambda$, or equivalently

$$\mathbf{q}_i^T A = \lambda_i \mathbf{q}_i^T \,,$$

for $1 \le i \le n$. Then

$$A^2 = (P\Lambda P^{-1})(P\Lambda P^{-1}) = P\Lambda^2 P^{-1}$$

and $A^k = P\Lambda^k P^{-1}$, so that

$$\exp(At) = \sum_{k=0}^{\infty} \frac{(At)^k}{k!} = P\left[\sum_{k=0}^{\infty} \frac{(\Lambda t)^k}{k!}\right] P^{-1}$$
$$= P \exp(\Lambda t) P^{-1}. \tag{10}$$

The main advantage of expression (10) is that, since Λ is diagonal, its exponential matrix is easy to compute and is also diagonal, i.e.

$$e^{\Lambda t} = \operatorname{diag} \left\{ \exp(\lambda_i t), 1 \le i \le n \right\}.$$

The identity (10) expresses $\exp(At)$ completely in terms of the eigenvalues and right and left eigenvectors of A. To see this, note that

can be rewritten as

$$e^{At} = \sum_{i=1}^{n} \mathbf{p}_i \mathbf{q}_i^T \exp(\lambda_i t) . \tag{11}$$

Example: Let

$$A = \left[\begin{array}{ccc} 1 & 0 & -2 \\ 0 & -1 & 0 \\ 6 & 0 & -6 \end{array} \right] .$$

Then

$$sI - A = \left[\begin{array}{ccc} s - 1 & 0 & 2 \\ 0 & s + 1 & 0 \\ -6 & 0 & s + 6 \end{array} \right]$$

and

$$a(s) = \det(sI - A) = (s+1)[(s-1)(s+6) + 12] = (s+1)(s+2)(s+3).$$

The eigenvector \mathbf{p}_1 corresponding to $\lambda_1 = -1$ is obtained by solving

$$(\lambda_1 I - A)\mathbf{p}_1 = \left[egin{array}{ccc} -2 & 0 & 2 \\ 0 & 0 & 0 \\ -6 & 0 & 5 \end{array} \right] \left[egin{array}{c} x_1 \\ x_2 \\ x_3 \end{array} \right] = \left[egin{array}{c} 0 \\ 0 \\ 0 \end{array} \right] \,.$$

This gives $x_1 = x_3 = 0$ with x_2 free, and since the scaling of \mathbf{p}_1 is arbitrary, we select

$$\mathbf{p}_1 = \left[\begin{array}{c} 0 \\ 1 \\ 0 \end{array} \right] \ .$$

Similarly the eigenvector \mathbf{p}_2 corresponding to $\lambda_2 = -2$ is obtained by solving

$$(\lambda_2 I - A)\mathbf{p}_2 = \begin{bmatrix} -3 & 0 & 2 \\ 0 & -1 & 0 \\ -6 & 0 & 4 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix},$$

which gives $x_2 = 0$ and $3x_1 = 2x_3$, so that we can select

$$\mathbf{p}_2 = \left[egin{array}{c} 2 \ 0 \ 3 \end{array}
ight] \, .$$

Finally, to find eigenvector \mathbf{p}_3 corresponding to $\lambda_3 = -3$, we solve

$$(\lambda_3 I - A)\mathbf{p}_3 = \left[egin{array}{ccc} -4 & 0 & 2 \ 0 & -2 & 0 \ -6 & 0 & 3 \end{array} \right] \left[egin{array}{c} x_1 \ x_2 \ x_3 \end{array} \right] = \left[egin{array}{c} 0 \ 0 \ 0 \end{array} \right] \, .$$

This gives $x_2 = 0$ and $x_3 = 2x_1$, so that we can choose

$$\mathbf{p}_3 = \left[\begin{array}{c} 1 \\ 0 \\ 2 \end{array} \right] .$$

Inverting the matrix

$$P = \left[egin{array}{ccc} \mathbf{p}_1 & \mathbf{p}_2 & \mathbf{p}_3 \end{array}
ight] \left[egin{array}{ccc} 0 & 2 & 1 \ 1 & 0 & 0 \ 0 & 3 & 2 \end{array}
ight]$$

gives

$$Q = P^{-1} = \left[\begin{array}{rrr} 0 & 1 & 0 \\ 2 & 0 & -1 \\ -3 & 0 & 2 \end{array} \right] ,$$

so that

$$\exp(At) = \begin{bmatrix} 0 & 2 & 1 \\ 1 & 0 & 0 \\ 0 & 3 & 2 \end{bmatrix} \begin{bmatrix} \exp(-t) & 0 & 0 \\ 0 & \exp(-2t) & 0 \\ 0 & 0 & \exp(-3t) \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 \\ 2 & 0 & -1 \\ -3 & 0 & 2 \end{bmatrix} \\
= \begin{bmatrix} 4\exp(-2t) - 3\exp(-3t) & 0 & 2[-\exp(-2t) + \exp(-3t)] \\ 0 & \exp(-t) & 0 \\ 6[\exp(-2t) - \exp(-3t)] & 0 & -3\exp(-2t) + 4\exp(-3t) \end{bmatrix}.$$

In the expression (11) for $\exp(At)$, although $\exp(At)$ is real, the eigenvalues λ_i and eigenvectors \mathbf{p}_i and \mathbf{q}_i may be *complex*.

Example: Let

$$A = \left[\begin{array}{rrr} -2 & 0 & -1 \\ 0 & -2 & 0 \\ 2 & 0 & 0 \end{array} \right] .$$

Then

$$sI - A = \left[\begin{array}{ccc} s+2 & 0 & 1 \\ 0 & s+2 & 0 \\ -2 & 0 & s \end{array} \right] ,$$

and

$$a(s) = \det sI - A = (s^2 + 2s + 2)(s + 2) = (s + 1 - j)(s + 1 + j)(s + 2)i$$
.

Thus A has two complex conjugate eigenvalues $\lambda_{\pm} = -1 \pm j$ and a real eigenvalue $\lambda_3 = -2$. The eigenvector \mathbf{p}_{+} coorresponding to λ_{+} is obtained by solving

$$(\lambda_+ I - A)\mathbf{p}_+ = \left[egin{array}{ccc} 1+j & 0 & 1 \ 0 & 1+j & 0 \ -2 & 0 & -1+j \end{array}
ight] \left[egin{array}{c} x_1 \ x_2 \ x_3 \end{array}
ight] = \left[egin{array}{c} 0 \ 0 \ 0 \end{array}
ight] \, ,$$

which gives $x_2 = 0$, $x_3 = -(1+j)x_1$, so that we can select

$$\mathbf{P}_{+} = \left[\begin{array}{c} 1 \\ 0 \\ -(1+j) \end{array} \right] .$$

By observing that $\lambda_- = \lambda_+^*$ and A is real, and taking the complex conjugate of the eigenvector equation for λ_+ we find that

$$\mathbf{p}_- = \mathbf{p}_+^* = \left[egin{array}{c} 1 \ 0 \ -(1-j) \end{array}
ight]$$

satisfies

$$(\lambda_- I - A)\mathbf{p}_- = \mathbf{0} ,$$

so it is the eigenvector corresponding to eigenvalue λ_{-} . Finally, the eigenvector corresponding to $\lambda_{3} = -2$ is obtained by solving

$$(\lambda_3 I - A)\mathbf{p}_3 = \left[egin{array}{ccc} 0 & 0 & 1 \ 0 & 0 & 0 \ -2 & 0 & -2 \end{array} \right] \left[egin{array}{c} x_1 \ x_2 \ x_3 \end{array} \right] = \left[egin{array}{c} 0 \ 0 \ 0 \end{array} \right] \, ,$$

which yields $x_1 = x_3 = 0$, so that we can select

$$\mathbf{p}_3 = \left[egin{array}{c} 0 \ 1 \ 0 \end{array}
ight] \; .$$

Then if

$$P = \left[egin{array}{cccc} 1 & 1 & 0 \ 0 & 0 & 1 \ -(1+j) & -(1-j) & 0 \end{array}
ight]$$

we find

$$Q = P^{-1} = \begin{bmatrix} (1+j)/2 & 0 & j/2 \\ (1-j)/2 & 0 & -j/2 \\ 0 & 1 & 0 \end{bmatrix}$$

and

$$\exp(At) = P \exp(\Lambda t)Q$$

with

$$\exp(\Lambda t) = \left[egin{array}{ccc} \exp(-t+jt) & 0 & 0 \ 0 & \exp(-t-jt) & 0 \ 0 & 0 & \exp(-2t) \end{array}
ight] \, .$$

Real form of the eigenvalue/eigenvector expansion of $\exp(At)$: To obtain expressions involving only real quantities, one can proceed as follows. First observe that since A is a real matrix, if $(\lambda_i, \mathbf{p}_i)$ is an eigenvalue/eigenvector pair for A, so is $(\lambda_i^*, \mathbf{p}_i^*)$. Assume now that A has r pairs of complex conjugate eigenvalues and q = n - 2r real eigenvalues. Then order the eigenvalues of A as $\{\lambda_1, \lambda_1^*, \dots, \lambda_k, \lambda_k^*, \dots, \lambda_r, \lambda_r^*, \lambda_{2r+1}, \dots, \lambda_n\}$. In the eigenvalue/eigenvector relation $A\mathbf{p}_k = \lambda_k \mathbf{p}_k$, the eigenvalue λ_k and eigenvector \mathbf{p}_k can be decomposed into their real and imaginary parts as

$$\lambda_k = a_k + jb_k$$
$$\mathbf{p}_k = \mathbf{p}_k^R + j\mathbf{p}_k^I.$$

This implies that

$$A\mathbf{p}_k^R = a_k \mathbf{p}_k^R - b_k \mathbf{p}_k^I$$

$$A\mathbf{p}_k^I = b_k \mathbf{p}_k^R + a_k \mathbf{p}_k^I$$

Then, if we consider the matrix

$$M = \begin{bmatrix} \mathbf{p}_1^R \mathbf{p}_1^I & \dots & \mathbf{p}_k^R \mathbf{p}_k^I & \dots & \mathbf{p}_r^R \mathbf{p}_r^I & \mathbf{p}_{2r+1} & \dots & \mathbf{p}_n \end{bmatrix}$$

M is real and

$$AM = MD$$

with

where

$$D_k \stackrel{ riangle}{=} \left[egin{array}{cc} a_k & b_k \ -b_k & a_k \end{array}
ight]$$

If $N = M^{-1}$, this implies

$$A = MDN \tag{12}$$

where D is block diagonal. It is constituted of $r \, 2 \times 2$ blocks corresponding to the r pairs $(\lambda_k, \, \lambda_k^*)$ of complex conjugate eigenvalues, as well as $n - 2r \, 1 \times 1$ blocks corresponding to the n - 2r real eigenvalues $\lambda_{2r+1}, \ldots, \lambda_n$ of A.

The expression (12) for A implies

$$\exp(At) = M \exp(Dt)N, \qquad (13)$$

with

and

$$\exp(D_k t) = \exp(a_k t) \begin{bmatrix} \cos(b_k t) & \sin(b_k t) \\ -\sin(b_k t) & \cos(b_k t) \end{bmatrix},$$

where the expression (13) involves only real matrices.

Example: It was found earlier that the matrix

$$A = \left[\begin{array}{rrr} -2 & 0 & -1 \\ 0 & -2 & 0 \\ 2 & 0 & 0 \end{array} \right]$$

has eigenvalues $\lambda_{\pm} = -1 \pm j$ and $\lambda_{3} = -2$. Its eigenvectors are

$$\mathbf{p}_{\pm} = \left[egin{array}{c} 1 \ 0 \ -(1 \pm j) \end{array}
ight] \quad ext{and} \quad \mathbf{p}_{3} = \left[egin{array}{c} 0 \ 1 \ 0 \end{array}
ight]$$

The vector \mathbf{p}_+ can be decomposed as $\mathbf{p}_+ = \mathbf{p}_R + j\mathbf{p}_I$ with

$$\mathbf{p}_R = \begin{bmatrix} 1 \\ 0 \\ -1 \end{bmatrix} \quad \text{and} \quad \mathbf{p}_I = \begin{bmatrix} 0 \\ 0 \\ -1 \end{bmatrix}$$

Let

$$M \; = \; \left[egin{array}{ccc} {f p}_R & {f p}_I & {f p}_3 \end{array}
ight] = \left[egin{array}{ccc} 1 & 0 & 0 \ 0 & 0 & 1 \ -1 & -1 & 0 \end{array}
ight]$$

$$N = M^{-1} = \left[egin{array}{ccc} 1 & 0 & 0 \ -1 & 0 & -1 \ 0 & 1 & 0 \end{array}
ight]$$

$$D = \begin{bmatrix} -1 & 1 & 0 \\ -1 & -1 & 0 \\ 0 & 0 & -2 \end{bmatrix}.$$

Then

$$\exp(At) = M \exp(Dt)N$$

$$= \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ -1 & -1 & 0 \end{bmatrix} \begin{bmatrix} \exp(-t)\cos t & \exp(-t)\sin t & 0 \\ -\exp(-t)\sin t & \exp(-t)\cos t & 0 \\ 0 & 0 & \exp(-2t) \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ -1 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix}.$$

Case 2: A is not diagonalizable. In this case A has a Jordan block of size 2 or larger, or equivalently it has a repeated eigenvalue which admits fewer independent eigenvectors than its multiplicity. In this case, we can find an invertible matrix T constituted of the eigenvectors and generalized eigenvectors of A such that

$$A = TJT^{-1} (14)$$

with

$$J = \begin{bmatrix} J_1 & & & 0 \\ & \ddots & & \\ & & J_i & \\ & & & \ddots & \\ 0 & & & J_\ell \end{bmatrix}$$

$$J_i = \left[egin{array}{cccc} \lambda_i & 1 & & 0 \ & \ddots & \ddots & \ & & \lambda_i & 1 \ 0 & & & \lambda_i \end{array}
ight] = n_i imes n_i ext{ Jordan block },$$

where $\sum_{i=1}^{\ell} n_i = n$. The matrix exponential $\exp(At)$ is then given by

$$\exp(At) = T \exp(Jt)T^{-1}, \qquad (15)$$

with

$$\exp(Jt) = \left[egin{array}{ccc} \exp(J_1t) & & & 0 \ & \ddots & & & \ & \exp(J_it) & & & \ & & \ddots & & \ & & & \ddots & & \ & & & \exp(J_\ell t) \end{array}
ight]$$

and

$$\exp(J_i t) = \exp(\lambda_i t) \left[egin{array}{cccc} 1 & t & & rac{t^{n_i-1}}{(n_i-1)!} \ & \ddots & \ddots & \ & & \ddots & t \ 0 & & 1 \end{array}
ight].$$

Example: Let

$$A = \left[egin{array}{cccc} 1 & 2 & 0 & 1 \ 0 & 1 & 0 & 0 \ 0 & -1 & 1 & 0 \ 0 & 0 & 0 & 1 \end{array}
ight] \,.$$

Then

$$sI_4 - A = \left[egin{array}{cccc} s-1 & -2 & 0 & -1 \ 0 & s-1 & 0 & 0 \ 0 & 1 & s-1 & 0 \ 0 & 0 & 0 & s-1 \end{array}
ight] \, ,$$

and using Laplace's expansion of $a(s) = \det sI - A$ with respect to the first column yields $a(s) = (s-1)^4$. The eigenvectors \mathbf{p} of A corresponding to $\lambda_1 = 1$ satisfy

$$(\lambda_1 I - A) \mathbf{p} = \left[egin{array}{cccc} 0 & -2 & 0 & -1 \ 0 & 0 & 0 & 0 \ 0 & 1 & 0 & 0 \ 0 & 0 & 0 & 0 \end{array}
ight] \left[egin{array}{c} x_1 \ x_2 \ x_3 \ x_4 \end{array}
ight] = \left[egin{array}{c} 0 \ 0 \ 0 \ 0 \end{array}
ight] \,.$$

This gives $x_2 = x_4 = 0$, so that A has the two eigenvectors

$$\mathbf{p}_1 = \left[egin{array}{c} 1 \ 0 \ 0 \ 0 \end{array}
ight] \qquad \mathbf{p}_2 = \left[egin{array}{c} 0 \ 0 \ 1 \ 0 \end{array}
ight] \;.$$

Since the number of eigenvectors (two) is less than the multiplicity (four) of λ_1 , A is not diagonalizable. Observing that $(\lambda_1 I - A)^2 = 0$, we can conclude that A will have 2 Jordan blocks of size 2×2 corresponding to $\lambda_1 = 1$. The generalized eigenvector \mathbf{g}_1 corresponding to \mathbf{p}_1 is given by

$$(A-\lambda_1 I)\mathbf{g}_1 = \left[egin{array}{cccc} 0 & 2 & 0 & 1 \ 0 & 0 & 0 & 0 \ 0 & -1 & 0 & 0 \ 0 & 0 & 0 & 0 \end{array}
ight] \left[egin{array}{c} x_1 \ x_2 \ x_3 \ x_4 \end{array}
ight] = \mathbf{p}_1 = \left[egin{array}{c} 1 \ 0 \ 0 \ 0 \ 0 \end{array}
ight] \; ,$$

which gives $x_2 = 0$ and $x_4 = 1$, so that we can select

$$\mathbf{g}_1 = \left[\begin{array}{c} 0 \\ 0 \\ 0 \\ 1 \end{array} \right] \ .$$

Similarly the generalized eigenvector \mathbf{g}_2 corresponding to \mathbf{p}_2 satisfies

$$(A-\lambda_1 I)\mathbf{g}_2 = \left[egin{array}{cccc} 0 & 2 & 0 & 1 \ 0 & 0 & 0 & 0 \ 0 & -1 & 0 & 0 \ 0 & 0 & 0 & 0 \end{array}
ight] \left[egin{array}{c} x_1 \ x_2 \ x_3 \ x_4 \end{array}
ight] = \mathbf{p}_2 = \left[egin{array}{c} 0 \ 0 \ 1 \ 0 \end{array}
ight] \,.$$

We obtain $x_2 = -1$ and $x_4 = 2$, so that we can select

$$\mathbf{g}_2 = \left[egin{array}{c} 0 \ -1 \ 0 \ 2 \end{array}
ight] \, .$$

Then if

$$T^{-1} = \left[egin{array}{cccc} 1 & 0 & 0 & 0 \ 0 & 2 & 0 & 1 \ 0 & 0 & 1 & 0 \ 0 & -1 & 0 & 0 \end{array}
ight]$$

$$J = \left[egin{array}{cccc} 1 & 1 & 0 & 0 \ 0 & 1 & 0 & 0 \ 0 & 0 & 1 & 1 \ 0 & 0 & 0 & 1 \end{array}
ight],$$

we have $A = TJT^{-1}$ and

$$\exp(At) = T\exp(Jt)T^{-1}$$

with

$$\exp(Jt) = \exp(t) \left[egin{array}{cccc} 1 & t & 0 & 0 \ 0 & 1 & 0 & 0 \ 0 & 0 & 1 & t \ 0 & 0 & 0 & 1 \end{array}
ight] \; .$$

Note that the Jordan decomposition $A = TJT^{-1}$ may be complex. There exists a real form of this decomposition, but it will not be needed in the remainder of this course.